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Observations of the Flow Patterns of a  
Two-Dimensional 4 per cent Thick  
Biconvex Aerofoil at  $M_0 = 1.40$  and  $1.63$

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# Observations of the Flow Patterns of Two-Dimensional 4 per cent Thick Biconvex Aerofoil at $M_0 = 1.40$ and $1.63$

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*Summary.*—Direct-shadow and schlieren photographs and pressure distributions of the flow past a two-dimensional 4 per cent thick biconvex aerofoil for a limited range of incidences at Mach numbers of 1.40 and 1.63 are presented.

Shock-induced boundary-layer separation at the trailing edge of the aerofoil was present at  $M_0 = 1.63$  with transition-free boundary layers but was absent up to 5 deg incidence at  $M_0 = 1.63$  with transition-fixed boundary layers and up to 6 deg incidence at  $M_0 = 1.40$  with transition-free boundary layers.

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1. *Introduction.*—In previous reports<sup>1,2</sup> investigations of the flow past a two-dimensional 4 per cent thick biconvex aerofoil at subsonic speeds have been made. Recently tests on the same aerofoil at two supersonic speeds, namely free-stream Mach numbers of 1.40 and 1.63, have been completed and the results are presented. In a subsequent paper a description will be given of the testing of this biconvex aerofoil in slotted-wall transonic liners and thus results will be available from low subsonic ( $M_0 = 0.40$ ) to moderate supersonic speeds for wide ranges of incidence.

The Reynolds numbers of the current tests, based on the aerofoil chord of 9 in., were approximately  $3.4 \times 10^6$  and  $3.5 \times 10^6$  at  $M_0 = 1.63$  and 1.40 respectively. Throughout this investigation, particular attention was given to the presence (or otherwise) of shock-induced boundary-layer separation at the trailing edge of the aerofoil<sup>3</sup>.

2. *Experimental Data.*—2.1. *Results at  $M_0 = 1.63$ .*—Previous experiments on this 4 per cent biconvex aerofoil in the 36-in.  $\times$  14-in. High-Speed Wind Tunnel at subsonic speeds<sup>1,2</sup> had shown that only minute differences existed between tests made with natural (or transition-free) boundary layers and those with transition-fixed boundary layers produced artificially by 5 per cent chord bands of carborundum at the leading edge of the aerofoil. Hence the first series of tests at  $M_0 = 1.63$  were made using transition-free boundary layers. When shock-induced boundary-layer separation at the trailing edge of the aerofoil was found to be present, a further series of tests were made with transition-fixed boundary layers and in these results no regions of separated flow were observed. Figs. 1 and 2 present the variation of normal-force coefficient  $C_N$  and pitching-moment coefficient  $C_m$  with incidence at  $M_0 = 1.63$ . Whilst distinct differences exist between the transition-free and the transition-fixed results, the pressure distributions given in Figs. 3 and 4 show that these differences are confined to the trailing-edge region. Schlieren photographs such as Fig. 5 confirm these remarks but the two distinct flow patterns are most clearly illustrated by the direct-shadow photographs presented in Fig. 6. With transition-fixed

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boundary layers there is no separation at the trailing edge but with transition-free boundary layers separation is present at  $\alpha = 3$  deg and increases in chordwise extent as the incidence is increased. At  $\alpha = 5$  deg the separated region extends from 0.95 chord to the trailing edge.

It is of interest to determine the flow deflection which will just cause flow separation. For a circular-arc biconvex aerofoil the trailing-edge semi-angle  $\beta$  is approximately given by

$$\tan \beta \simeq 2 \times \frac{\text{aerofoil thickness}}{\text{aerofoil chord}}.$$

Hence for this 4 per cent biconvex aerofoil

$$\beta \simeq \tan^{-1} \frac{8}{100} \simeq 4^{\circ}36'.$$

Thus the flow deflection angle  $\theta$  and the incidence of the aerofoil  $\alpha$  are connected by the relation

$$\theta \simeq \alpha + 4^{\circ}36' \simeq \alpha + 4\frac{1}{2}^{\circ}.$$

Ref. 3 reports an investigation of the interaction between shock waves and boundary layers at the trailing edge of a double-wedge aerofoil at  $M_0 = 1.60$ . For the transition-free case, where the boundary layer was laminar at separation but turbulent at reattachment to wake, laminar separation was present when  $\theta = 5$  deg. For the transition-fixed or turbulent case, turbulent separation was present when  $\theta = 12\frac{1}{2}$  deg. Hence, in the present investigation, shock-induced boundary-layer separation might be expected to be observed at  $\alpha = 1$  deg in the transition-free case and  $\alpha = 8$  deg in the transition-fixed case. However, at  $M_0 = 1.63$  the tunnel became choked above  $\alpha = 5$  deg and investigation of the turbulent case was impossible. The transition-free results show separated regions present from  $\alpha = 3$  deg onwards, that is,  $\theta = 7\frac{1}{2}$  deg onwards.

The direct determination of the occurrence of separation by divergence of the trailing-edge pressure<sup>5</sup> was attempted in Fig. 7, but more consistent results were obtained in Fig. 8 where the divergence of the pressure at 0.97 chord was considered (this may be attributed to the diffuse nature of the trailing-edge shock wave in the supersonic case). From these figures it may be deduced that laminar separation is present at  $\theta = 7$  deg. The upstream influence of the shock-wave boundary-layer interaction was calculated using the methods of Refs. 3 and 4. When  $\alpha = 5$  deg, the extent of this influence is given by  $d/\delta_0^* = 32$ , where

$d$  = distance from trailing edge of beginning of separated region

$\delta_0^*$  = calculated boundary-layer displacement thickness at beginning of separated region.

This result is not inconsistent with the findings of Refs. 3 and 4. However, it should be noted that the Reynolds numbers of the present tests are an order of magnitude greater than for the previous tests<sup>3,4</sup>.

*2.2. Results at  $M_0 = 1.40$ .*—The variation of pitching-moment coefficient and normal-force coefficient with incidence at  $M_0 = 1.40$  is given in Figs. 1 and 2, and specimen pressure distributions are presented in Fig. 9.

Schlieren and direct-shadow photographs of the flow appear as Figs. 10 and 11 respectively.

These tests were made with transition-free boundary layers and no trace of shock-induced boundary-layer separation was observed in the photographs or could be deduced from the pressure distributions (incidences above  $\alpha = 6$  deg were not obtainable since the tunnel became choked). It is concluded that at  $M_0 = 1.40$  the boundary layers are of a naturally transitional type and are not strictly laminar at the beginning of the shock-wave boundary-layer interaction region.

*3. Conclusions.*—Shock-induced boundary-layer separation at the trailing edge of a 4 per cent thick biconvex aerofoil was present at  $M_0 = 1.63$  with transition-free boundary layers but was absent at  $M_0 = 1.63$  with transition-fixed boundary layers and at  $M_0 = 1.40$  with transition-free boundary layers. The extent of the separated region confirmed the results given in Ref. 3.

*Acknowledgements.*—Mr. P. J. Peggs assisted in the experimental work and Mrs. N. A. North in the data reduction.

## NOTATION

$M_0$	Tunnel free-stream Mach number
$M$	Local Mach number
$\theta$	Flow deflection angle
$\alpha$	Incidence of aerofoil
$\beta$	Trailing-edge semi-angle of aerofoil
$c$	Aerofoil chord
$d$	Length of separated flow region
$\delta_0^*$	Calculated boundary-layer displacement thickness at beginning of separated region
$C_N$	Normal-force coefficient
$C_m$	Pitching-moment coefficient (about 0.25 chord)
$p$	Local static pressure
$H_0$	Stagnation pressure
$x$	Distance measured from leading edge of aerofoil.

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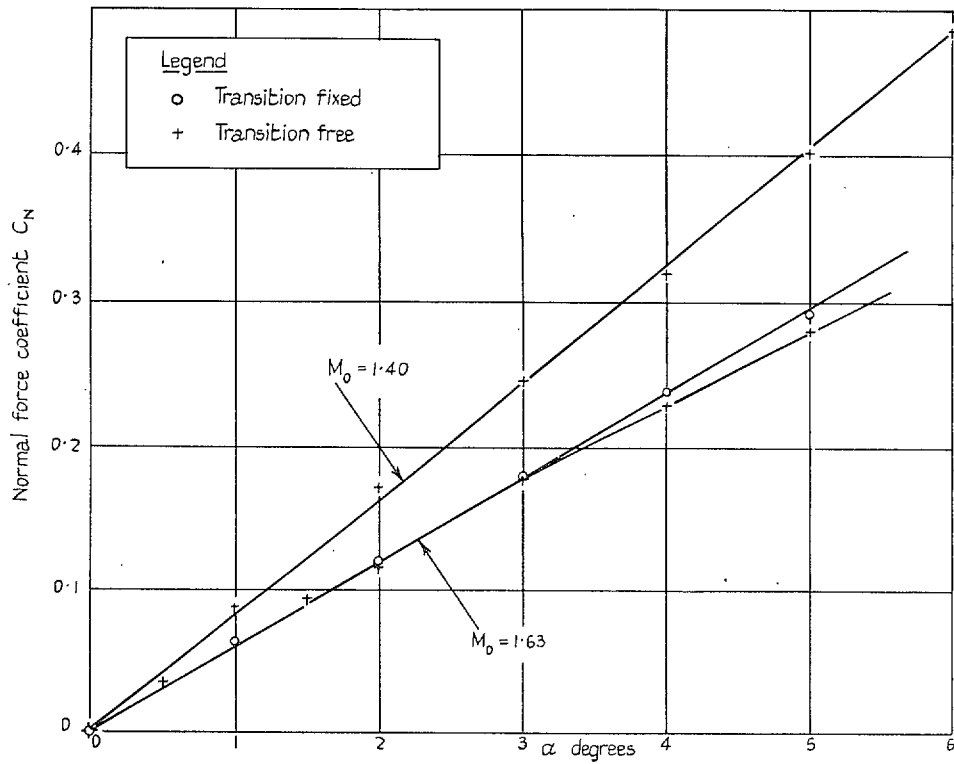


FIG. 1. Variation of normal-force coefficient  $C_N$  with incidence for 4 per cent biconvex aerofoil at  $M_0 = 1.63$  and  $M_0 = 1.40$ .

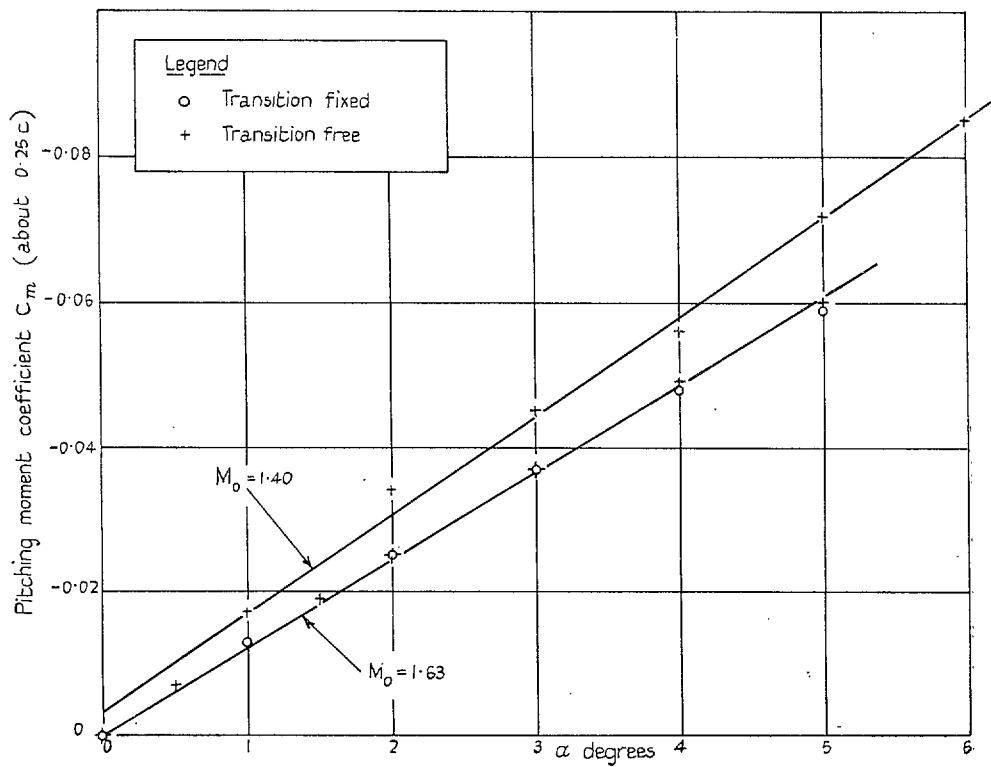
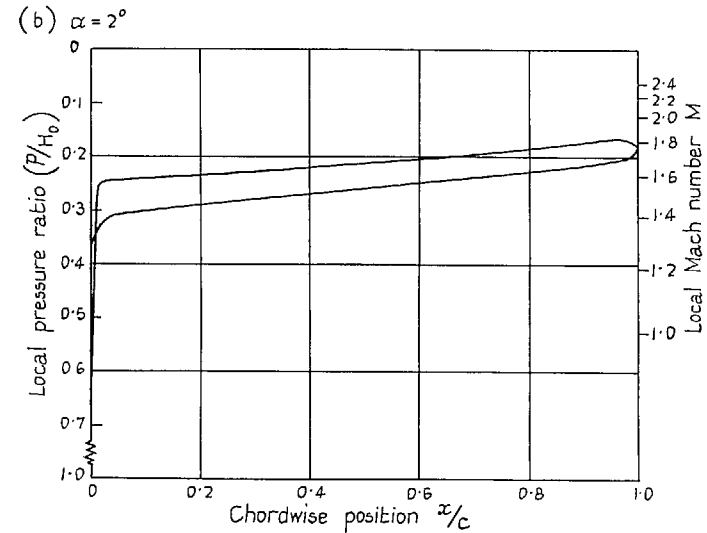
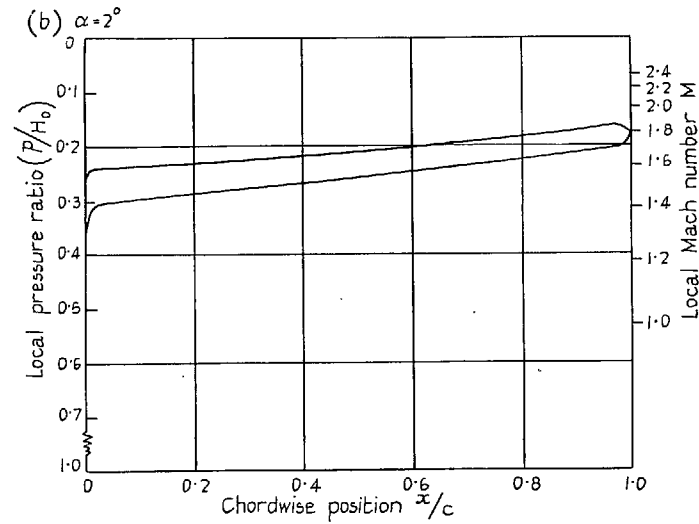
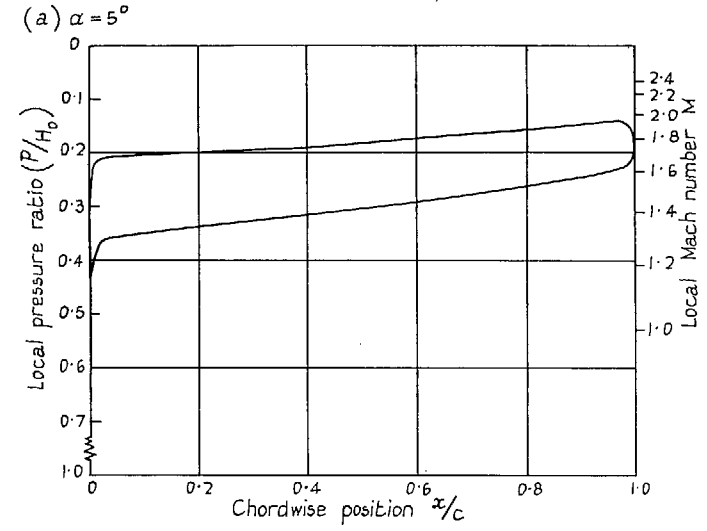
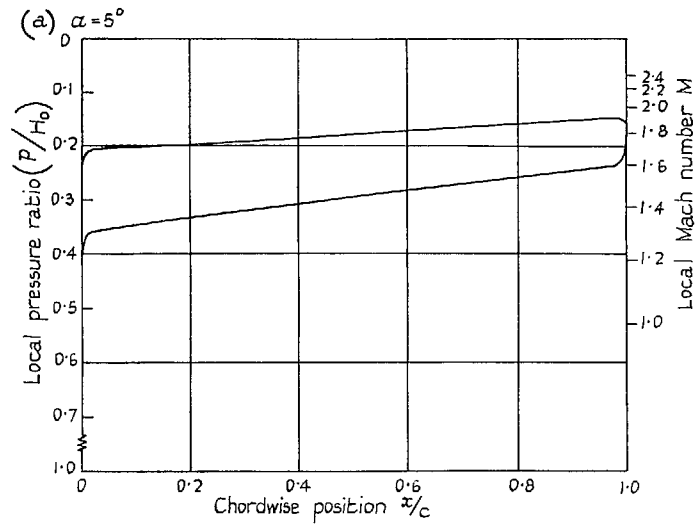
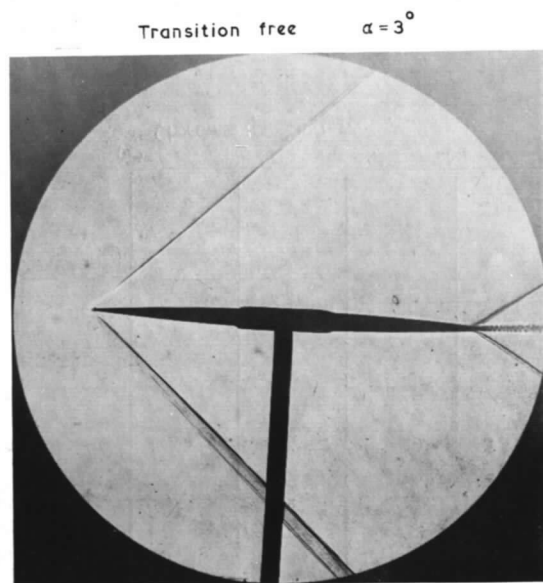
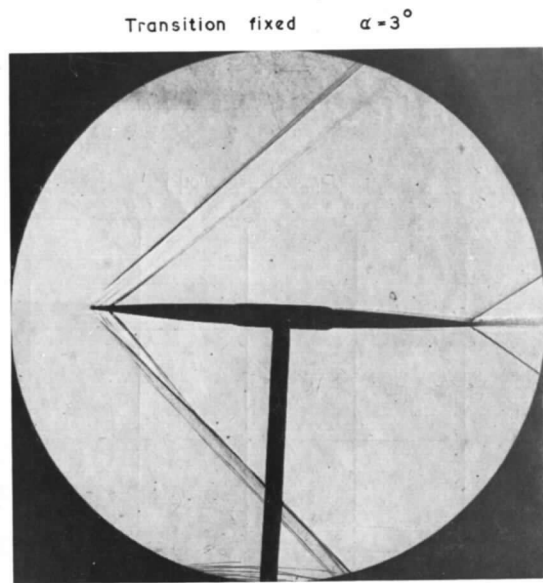


FIG. 2. Variation of pitching-moment coefficient  $C_m$  with incidence for 4 per cent biconvex aerofoil at  $M_0 = 1.63$  and  $M_0 = 1.40$ .



FIGS. 3a and 3b. Averaged pressure distributions for a 4 per cent biconvex aerofoil at  $M_0 = 1.63$  with 'transition-free' boundary layers.

FIGS. 4a and 4b. Averaged pressure distributions for a 4 per cent biconvex aerofoil at  $M_0 = 1.63$  with 'transition-fixed' boundary layers.



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FIG. 5. Schlieren photographs of the flow past a 4 per cent thick biconvex aerofoil at  $M_0 = 1.63$ .

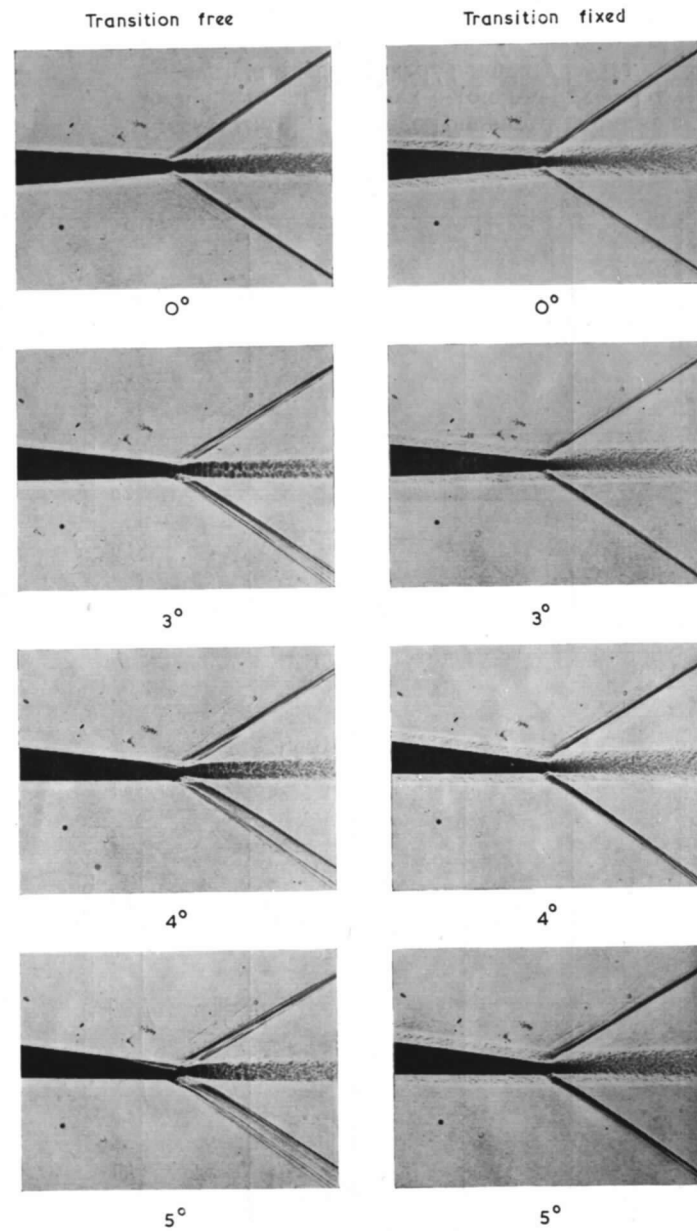


FIG. 6. Direct-shadow photographs of the flow past a 4 per cent thick biconvex aerofoil at  $M_0 = 1.63$ .

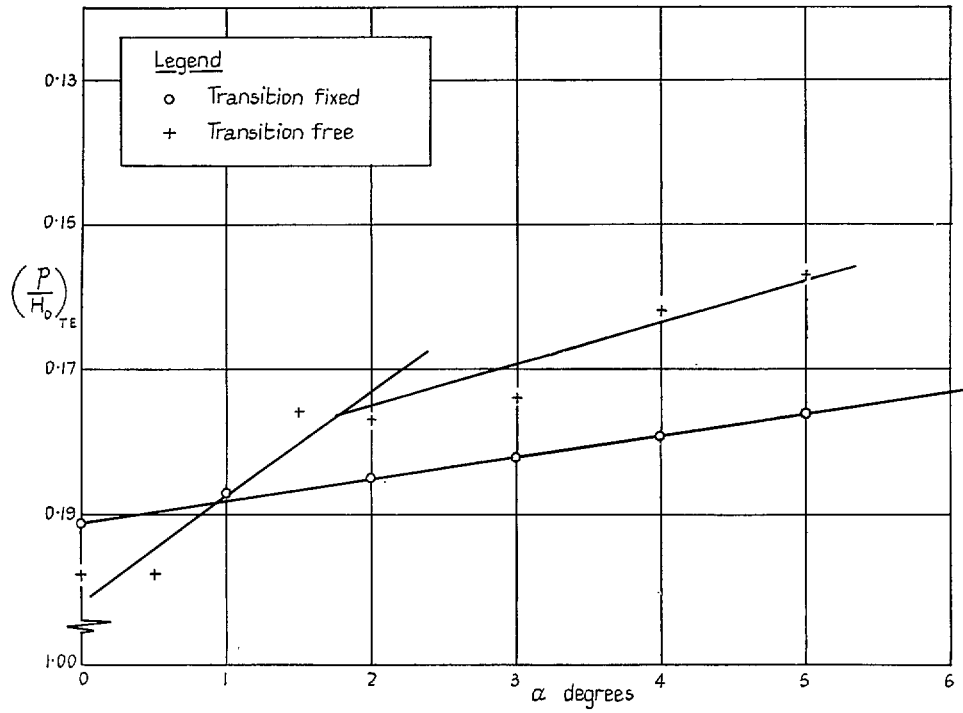


FIG. 7. Variation of the pressure ratio ( $p/H_0$ ) at the trailing edge with incidence for 4 per cent biconvex aerofoil at  $M_0 = 1.63$ .

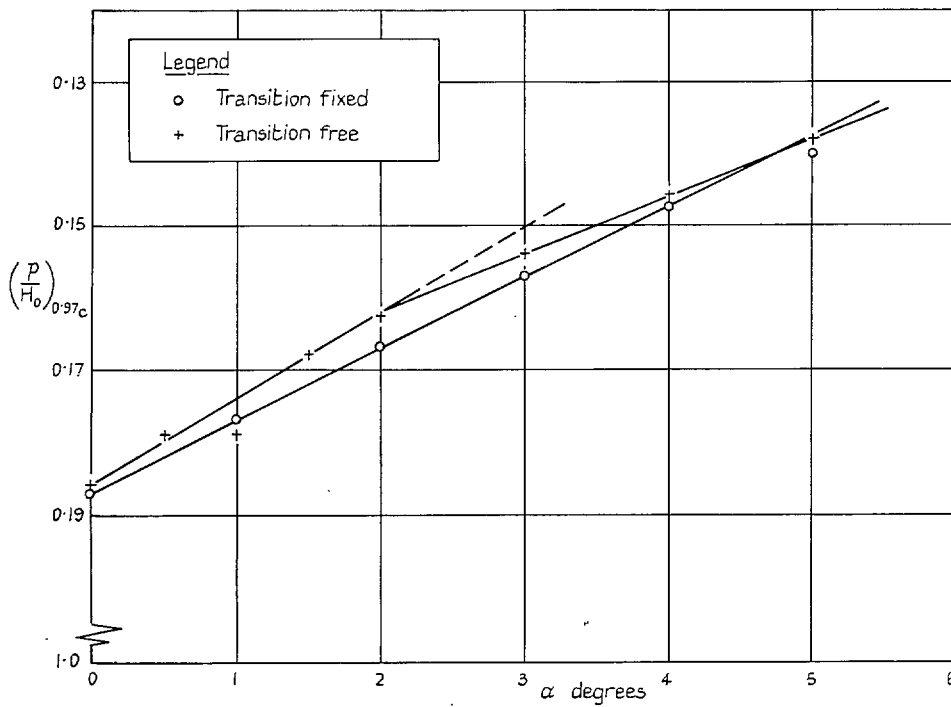
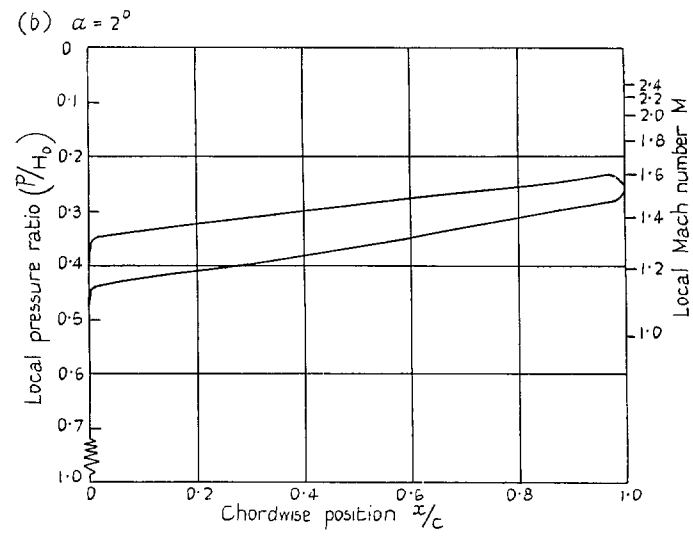
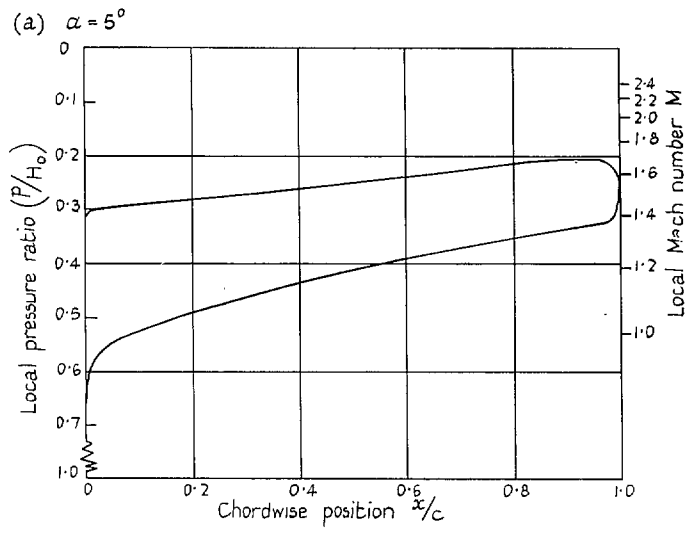


FIG. 8. Variation of the pressure ratio ( $p/H_0$ ) at  $0.97c$  with incidence for 4 per cent biconvex aerofoil at  $M_0 = 1.63$ .





FIGS. 9a and 9b. Averaged pressure distributions for a 4 per cent biconvex aerofoil at  $M_0 = 1.40$  with 'transition-free' boundary layers.

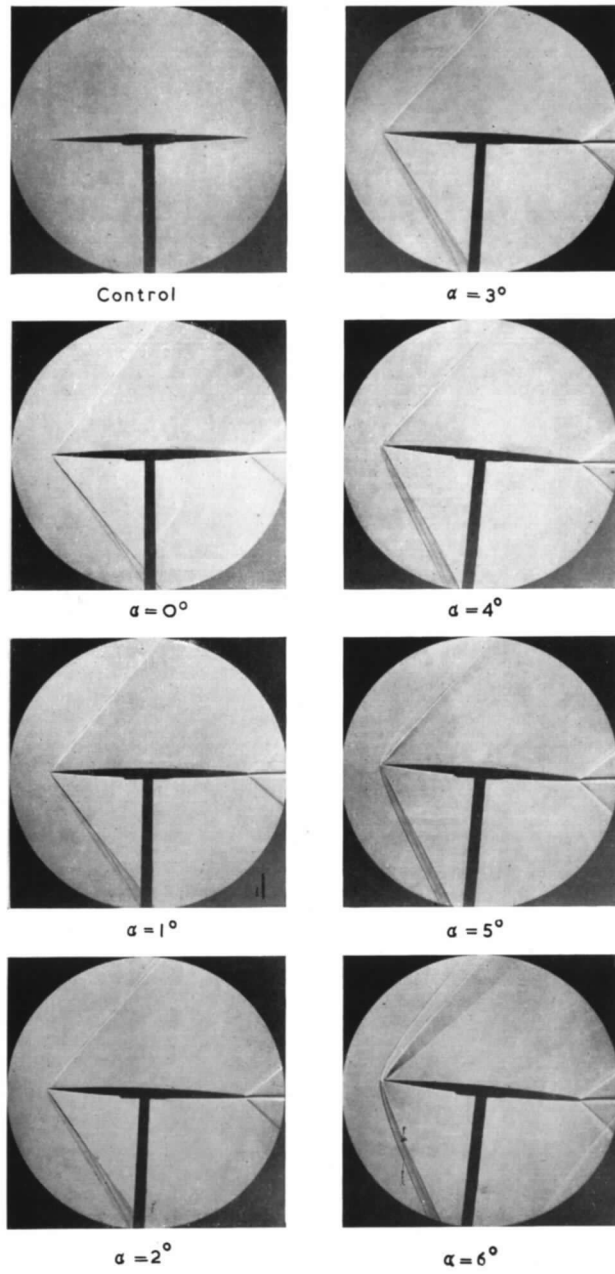


FIG. 10. Schlieren photographs of the flow past a 4 per cent thick biconvex aerofoil at  $M_0 = 1.40$ .

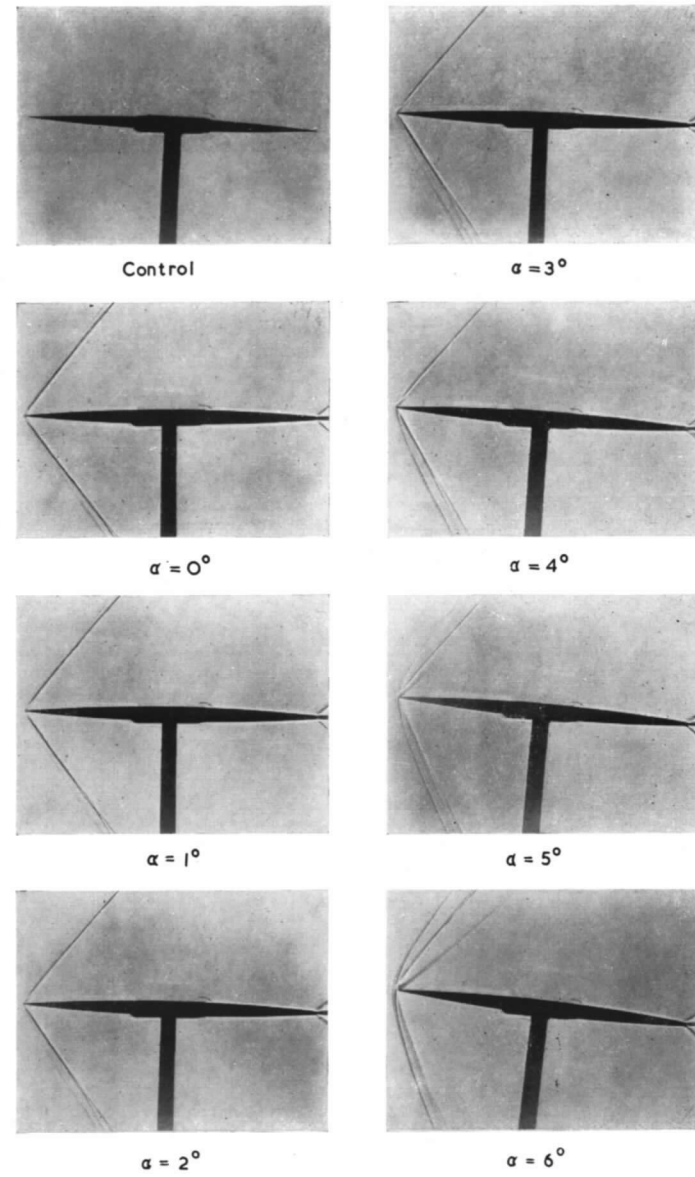


FIG. 11. Direct-shadow photographs of the flow past a 4 per cent thick biconvex aerofoil at  $M_0 = 1.40$ .

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